## Invisibility cloaking and cosmology

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Invisibility cloaking was used to find a counterexample for the inverse conductivity problem in Greenleaf-L.-Uhlmann 2003. The example was based on a blow-up map  $F: B(2)-\{0\} \rightarrow B(2)-\overline{B}(1)$ .

Using a similar map, Pendry, Schurig and Smith (Science 2006) suggested invisibility cloaking for Maxwell's equations, see also Leonhardt (Science 2006).



Figure: Implementation of an invisibility cloak for 4 cm waves build using metamaterials by Schurig et al (Science 2006).

# Cosmic Microwave Background measurements

The Friedmann-Robertson-Walker model of the Universe is the space-time  $M \times (0, \infty)$  with the metric

$$g=-dt^2+a(t)ig(\sum_{j,k=1}^3g_{jk}(x)dx^jdx^kig),\quad (x,t)\in M imes(0,\infty).$$

Here, M can be a compact or an unbounded manifold,  $\dim(M) = 3$ . In 2008, J-P. Luminet suggested that the cosmic microwave background in our Universe (the WMAP data) is compatible with this model, where M is Poincare dodecahedral space  $\mathbb{S}^3/\Gamma$ . Here,  $\Gamma$  is the binary icosahedral group.





Pictures by Jean-Pierre Luminet.

The right picture shows the Poincare dodecahedral space  $\mathbb{S}^3/\Gamma$ .

Next we consider a metamaterial device that can simulate waves in the space-time  $M \times (0, \infty)$  with a static metric  $g = -dt^2 + g_M(x)$ , Here, M is a 3-dimensional manifold with a Riemannian metric  $g_M$ . We proved an extension of the Lickorish-Wallace surgery theorem

### Proposition

Let M be a smooth, closed, oriented, and connected Riemannian manifold of dimension n. When n = 3, there exists a subset  $L \subset M$ , having capacity zero and Hausdorff dimension n - 2, such that there is a  $C^{\infty}$ -smooth embedding  $F : M - L \rightarrow \mathbb{R}^3$ . When n = 2, this result is not true if the genus of M is non-zero.

### Theorem (Cosmological cloaking)

Let (M,g) be a smooth, closed, oriented, and connected Riemannian manifold of dimension 3 and  $k^2 \notin \sigma(-\Delta_g)$ . Let  $\widetilde{M} = F(M-L) \subset \mathbb{R}^3$  and  $\widetilde{g} = F_*g$ . Also, let  $V \subset M-L$  and  $\widetilde{V} = F(V) \subset \widetilde{M}$ . Then  $\Delta_{\widetilde{g}} : H^2(\widetilde{M}, \widetilde{g}) \to L^2(\widetilde{M}, \widetilde{g})$  is self-adjoint and the source-to-solution map on (M, g),

$$L_V(s) = u|_V, \quad (\Delta_g + k^2)u = s, \quad s \in C_0^\infty(V),$$

coincides with the source-to-solution map  $L_{\widetilde{V}}$  on  $(\widetilde{M}, \widetilde{g})$ . A similar result holds for the approximative cloaking.

# A cosmological metamaterial device



Consider the interior of a twisted torus  $T \subset \mathbb{R}^3$  and remove several knotted, twisted toruses  $K_1, \ldots, K_J$ . This defines the domain  $\widetilde{M} = T - (\bigcup_{j=1}^J K_j)$ . Make gratings on the boundary  $\partial \widetilde{M}$  and cover it with the Schurig-type metamaterial. This construction gives a metamaterial device that simulates waves on the 3-manifold M.