

4. Quantum String

use always harmonic gauge, $h_{ab} = \begin{pmatrix} 0 & \sigma^+ \\ 0 & 0 \end{pmatrix}$ or $\begin{pmatrix} \sigma^+ & \sigma^- \\ 0 & -\frac{1}{2} \end{pmatrix}$

String dofs are

$$X^\mu(\tau, \sigma) \quad L = \frac{T}{2} \int d\sigma (\dot{X}^2 - X'^2) \quad P^\mu = \frac{\partial L}{\partial \dot{X}_\mu} = T \dot{X}^\mu$$

$$\mathcal{H} = P_\mu \frac{\partial L}{\partial \dot{X}_\mu} - L = \frac{T}{2} (\dot{X}^2 + X'^2)$$

or in the expansion

$$\underbrace{X^\mu}_{\text{real}} \quad \underbrace{P^\nu}_{\text{real}} \quad \underbrace{\alpha_m^\mu}_{\text{real}} \quad \bar{\alpha}_m^\mu$$

$$\alpha_m^\mu = \alpha_{-m}^\mu$$

Write commutators, generalising $[X, P^\nu] = i\eta^\nu$: $(-+++..)$

$$[X^\mu(\tau, \sigma), P^\nu(\tau, \sigma)] = T [X^\mu(\tau, \sigma), \dot{X}^\nu(\tau, \sigma)] = i \delta(\sigma - \sigma') \eta^{\mu\nu}$$

Insert the expansions here:

$[X^\mu, P^\nu] = i\eta^{\mu\nu}$	$[\alpha_m^\mu, \alpha_n^\nu] = \underset{\uparrow}{m} S_{m, -n} \eta^{\mu\nu}$	same for $\bar{\alpha}^\mu$, $[\alpha, \bar{\alpha}] = 0$
$\alpha_m \cdot \alpha_n - \alpha_n \cdot \alpha_m = m n S_{m, -n}$		

Harmonic oscillator: $[a, a^\dagger] = 1 \rightarrow [a_k, a_l^\dagger] = \delta_{kl}$

$$(aa^\dagger - a^\dagger a)|0\rangle = |0\rangle \Rightarrow aa^\dagger|0\rangle = |0\rangle \quad a|0\rangle = 0$$

$$N = a^\dagger a = aa^\dagger - 1 \quad \text{if } a^\dagger |m\rangle = m|m\rangle \text{ then}$$

$$N|a|m\rangle = (aN + Na - aN)|m\rangle = (aN - a)|m\rangle = (m-1)a|m\rangle$$

$$\underbrace{a^\dagger a a - a a^\dagger a}_{= a^\dagger a - 1}$$

$$N a^\dagger |m\rangle = \dots = (m+1)a^\dagger |m\rangle$$

$$\text{Normalised states: } |m\rangle = \frac{1}{\sqrt{m!}} (a^\dagger)^m |0\rangle \quad \begin{cases} a^\dagger |m\rangle = \sqrt{m+1} |m+1\rangle \\ a |m\rangle = \sqrt{m} |m-1\rangle \end{cases}$$

Coherent states: $a|z\rangle = z|z\rangle$
Solving this, $a = \frac{1}{\sqrt{2}} \left(\frac{x}{L} + i \frac{p_0}{\hbar} \right)$, gives $|z\rangle = e^{z a^\dagger - \frac{1}{2}|z|^2} |0\rangle$

$$\Psi_z(x) = \frac{1}{\sqrt{L\pi}} \exp \left[-\frac{1}{2L^2} (x - x_0)^2 + i \frac{p_0}{\hbar} x \right] \quad z = \frac{1}{\sqrt{2}} \left(\frac{x_0}{L} + i \frac{p_0}{\hbar} \right) \quad L^2 = \frac{\hbar}{mc\omega}$$

minimum
Uncertainty
state

$$\underbrace{[\alpha_{-m}^\mu, \alpha_m^\nu]} = -m \delta_{m,n} \gamma^{\mu\nu}$$

$$\alpha_{-m}^\mu \alpha_m^\nu - \alpha_m^\nu \alpha_{-m}^\mu = -m \gamma^{\mu\nu}$$

$$= \alpha_m^{\mu+} \quad \Rightarrow \underbrace{\left[\frac{1}{\sqrt{m}} \alpha_m^\nu, \frac{1}{\sqrt{m}} \alpha_m^{\mu+} \right]} = \delta_{mn} \gamma^{\mu\nu}$$

$$\alpha_m^\nu \equiv \alpha_m^{\mu+}$$

$$\Rightarrow [\alpha_m^\nu, \alpha_m^{\mu+}] = \delta_{mn} \gamma^{\mu\nu} \quad \gamma = (- + + + \dots)$$

↑
wrong sign!!

$\alpha_m^{\mu+}|0\rangle$ = state with osc. mode μ, m .

$$|\alpha_m^{0+}|0\rangle| = \langle 0 | \underbrace{\alpha_m^0 \alpha_m^{0+}}_{\alpha_m^{0+} \alpha_m^0} |0\rangle = -1 ! \quad \text{ghost state}$$

Including x^μ, p^μ : $|0\rangle \rightarrow |p\rangle \quad \hat{p}^\mu |p\rangle = p^\mu |p\rangle$

states: $|p\rangle \quad \alpha_{-1}^\mu, |p\rangle \quad \alpha_{-1}^\mu, \alpha_{-1}^\nu, \alpha_{-2}^\lambda |p\rangle \dots$ etc. Too many states; try to select physical ones by $L_k |p_{\text{phys}}\rangle = 0$ (cannot have $L_k = 0$!).

Operators: class \rightarrow quantum

how does one order non-commuting α 's?

$$L_k = \frac{1}{2} \sum_{n=-\infty}^{\infty} \alpha_{k-n}^\mu \alpha_{np} = \frac{1}{2} \left\{ \underbrace{\alpha_k^\mu \alpha_{0p}}_{\text{all commute for } k \neq 0! \text{ no problem!}} + \underbrace{\alpha_{k-1}^\mu \alpha_{1p}} + \underbrace{\alpha_{k+1}^\mu \alpha_{-1p}} + \dots \right.$$

$$L_0 = \frac{1}{2} \alpha_0^\mu \alpha_{0p} + \frac{1}{2} \sum_{m=1}^{\infty} \alpha_{-m}^\mu \alpha_{mp} + \frac{1}{2} \sum_{m=-1}^{-\infty} \alpha_{-m}^\mu \alpha_m^\nu$$

$$= \frac{1}{2} \left[\alpha_0 \cdot \alpha_0 + \underbrace{\alpha_{-1} \cdot \alpha_1}_{\text{normal order these}} + \underbrace{\alpha_1 \cdot \alpha_{-1}}_{d_{m<0} \cdot d_{m>0}} + \underbrace{\alpha_2 \cdot \alpha_2}_{\alpha_{-2} \cdot \alpha_2 + 2d} + \underbrace{\alpha_2 \cdot \alpha_{-2}}_{\dots} + \dots \right]$$

$$\alpha_m^\mu \cdot \alpha_{-m}^\nu = \alpha_{-m}^\mu \cdot \alpha_m^\nu + m d$$

$$= \frac{1}{2} \alpha_0 \cdot \alpha_0 + \sum_{m=1}^{\infty} \alpha_{-m} \cdot \alpha_m + \frac{1}{2} d \underbrace{(1+2+3+\dots)}_{\zeta(-1) = -\frac{1}{12}}$$

$$= L_0 - a$$

The correct one is
 $\frac{d-2}{24}$ (only $d-2$ trans.
physical dofs)

and this must be = 1
 $\Rightarrow d=26$ to have $m=0$
spin 1 states

$$\Rightarrow H(\text{closed}) = L_0 + \bar{L}_0 - 2a \quad (\text{p. 35})$$

$P = \bar{L}_0 - L_0$ generates translations in σ

$$\stackrel{\text{Cor R}}{\lambda a + \mu b}$$

Algebra $\hat{=}$ vector space with product $a \cdot b$

Grassmann: needed for superstrings!

$$c_i c_i = 0$$

m -dim vector space, basis c_1, \dots, c_m , product $c_i c_j = -c_j c_i$

Elements: 1 c_i $c_i c_j$ $c_i c_j c_k \dots$ $c_1 c_2 \dots c_m$

$$1 + \binom{m}{1} + \binom{m}{2} + \binom{m}{3} + \dots + \binom{m}{m} = (1+1)^m = 2^m$$

elements

Derivative: $\frac{\partial c_i}{\partial c_j} = \delta_{ij}$ $\frac{\partial}{\partial c_i} c_3 c_1 c_4 = -c_3 c_4$ etc

Integral: $\int dc_i = 0$ $\int dc_i c_j = - \int c_j dc_i = \delta_{ij}$

$$\int dc f(c) = \int dc [f(0) + c f'(0) + \dots] = \frac{\partial f}{\partial c} \quad (!)$$

$$\int d\bar{\psi}_1 d\bar{\psi}_2 d\psi_1 d\psi_2 = e^{\sum \bar{\psi}_i A_{ij} \psi_j} = A_{11} A_{22} - A_{12} A_{21} = \det A$$

$m=4$, 16 elements

$$\int d\bar{\psi} d\psi e^{-\bar{\psi} A \psi + \bar{\psi} \eta + \bar{\eta} \psi} = \det A \ e^{\bar{\eta} A^{-1} \eta} \quad \psi = (\psi_1, \dots, \psi_m) \text{ etc}$$

Lie: $[T_a, T_b] = i f_{abc} T_c$ $T_c^+ = T_c$, $U = e^{i \theta_a T_a}$ $= (U^+)^{-1}$ $a=1, \dots, \text{finite}$

Graded Lie: $[B, B] = B$ $[B, F] = F$ $[F, F] = B$

Virasoro $[L_m, L_n] = (m-n)L_{m+n} + \underbrace{\frac{c}{12} m(m^2-1)}_{\substack{\text{central extension, } c=\text{const} \\ \text{commutes with all } L_n}} S_{m,-n}$

Ex: Show that the Jacobi identity

$$[[L_m, L_n], L_p] + [[L_m, L_p], L_n] + [[L_p, L_m], L_n] = 0$$

is satisfied by the above form

Kac-Moody: $[T_m^a, T_n^b] = i f^{abc} T_{m+n}^c + \underbrace{k \cdot m}_{\substack{\text{commutes with all } T_m^a}} \delta^{ab} \delta_{m,-n}$

Now work out

$$[L_k, L_p] = \frac{1}{4} \sum_{m,n}^{\infty} [\alpha_{k-m}^r \alpha_m^v, \alpha_{p-m}^s \alpha_m^t] \gamma_{pr} \gamma_{st}$$

$$\begin{aligned} [ab, cd] &= abcd - cdab \\ &\quad \text{note that } [a, c] \text{ etc commute: } [\alpha_m^r, \alpha_m^v] = m \delta_{m,m} \gamma^{rv} \\ &= ac[b, d] + ad[b, c] + cb[a, d] + db[a, c] \\ &= \frac{1}{4} \sum_{m,n} \left[\alpha_{k-m}^r \alpha_{p-m}^s m \gamma^{rs} \delta_{m,-m} + \alpha_{k-m}^r \alpha_m^s m \gamma^{rs} \delta_{m,p-m} \right. \\ &\quad \left. + \alpha_{p-m}^s \alpha_m^r (k-m) \gamma^{rs} \delta_{k-m,-m} + \alpha_m^s \alpha_m^r (k-m) \gamma^{rs} \delta_{k-m, m-p} \right] \gamma_{pr} \gamma_{st} \\ &= \frac{1}{2} \sum_m \left[m \alpha_{k-m} \cdot \alpha_{p+m} + (k-m) \alpha_{p+k-m} \cdot \alpha_m \right] = \sum (k-m) \alpha_{p+k-m} \cdot \alpha_m \\ &\quad m=k-m' \quad = \frac{1}{2} [\text{same} - \text{same}(k+p)] = (k-p) L_{k+p} \end{aligned}$$

$$\Rightarrow [L_k, L_p] = (k-p) L_{k+p} \quad \text{if } p+k \neq 0$$

$$[L_k, L_{-k}] = \frac{1}{2} \sum_{-n}^{\infty} [m \alpha_{k-n} \cdot \alpha_{-k+n} + (k-n) \alpha_{-n} \cdot \alpha_n]$$

again the ordering problem since (see p. 43)

$$\alpha_m \cdot \alpha_{-m} - \alpha_{-m} \cdot \alpha_m = m \alpha_m !$$

$$\Rightarrow [L_k, L_p] = (k-p) L_{k+p} + \frac{d}{12} k(k^2-1) \delta_{k+p, 0}$$

Constraints $T_{++} = 0$ ($\tilde{L}_k = 0$) $T_{--} = 0$ ($L_{-k} = 0$)

can be expressed as

$$L_k | \text{phys} \rangle = 0 \quad \tilde{L}_k | \text{phys} \rangle = 0 \quad k > 0$$

$$(L_0 - a) | \text{phys} \rangle = 0 \quad (\tilde{L}_0 - a) | \text{phys} \rangle = 0 \quad k = 0 \quad \dots$$

but due to central term

$$\langle \text{ph} | L_k L_{-k} - \underbrace{L_{-k} L_k}_{=0} | \text{ph} \rangle = 2k \langle 0 | L_0 | 0 \rangle + (\text{const} \neq 0) \neq 0$$

$$\langle \text{ph} | L_k \rightarrow L_k^+ | \text{ph} \rangle = L_{-k}^- | \text{ph} \rangle \neq 0$$

Working out the central term:

Try $k=2$:

$$\left\{ \begin{array}{l} L_2 = \frac{1}{2} \sum \alpha_{2-m} \cdot \alpha_m = \frac{1}{2} \left[\alpha_4 \cdot \alpha_{-2} + \overbrace{\alpha_3 \cdot \alpha_{-1} + \alpha_2 \cdot \alpha_0 + \alpha_1 \cdot \alpha_1 + \alpha_0 \cdot \alpha_2 + \alpha_{-1} \cdot \alpha_3 + \dots}^{\text{central term}} \right] \\ = \frac{1}{2} \alpha_1 \cdot \alpha_1 + [\alpha_0 \cdot \alpha_2 + \alpha_{-1} \cdot \alpha_3 + \alpha_{-2} \cdot \alpha_4 + \dots] \\ L_{-2} = \frac{1}{2} \alpha_{-1} \cdot \alpha_{-1} + [\alpha_{-2} \cdot \alpha_0 + \alpha_{-3} \cdot \alpha_1 + \alpha_{-4} \cdot \alpha_2 + \dots] \end{array} \right.$$

$$[L_2, L_{-2}] = 4L_0 + \frac{d}{2}$$

The interesting part is $\alpha_m \cdot \alpha_m - \alpha_{-m} \cdot \alpha_m = m d \delta_{m+m,0}$

$$\begin{aligned} & \frac{1}{4} [\alpha_1 \cdot \alpha_1, \alpha_{-1} \cdot \alpha_{-1}] & \alpha_1 \cdot \alpha_1 - \alpha_{-1} \cdot \alpha_1 = d \\ & = \frac{1}{4} \left\{ [\alpha_1 \cdot \alpha_1, \alpha_{-1}^M] \alpha_{-1}^M + \alpha_{-1}^M [\alpha_1 \cdot \alpha_1, \alpha_{-1}^M] \right\} \\ & = [\alpha_1^\alpha \alpha_1, \alpha_{-1}^M] = \alpha_1^\alpha [\alpha_1, \alpha_{-1}^M] + [\alpha_1, \alpha_{-1}^M] \alpha_1^\alpha \\ & \quad \hat{=} 2 \underbrace{\alpha_1^\alpha [\alpha_1, \alpha_{-1}^M]}_{\eta_\alpha^M} = 2 \alpha_1^M \\ & = \frac{1}{2} (\underbrace{\alpha_1 \cdot \alpha_{-1}}_d + \underbrace{\alpha_{-1} \cdot \alpha_1}_{\text{normal ordered}}) & \eta_\alpha^M \\ & = \frac{d}{2} + \alpha_{-1} \cdot \alpha_1 \end{aligned}$$

this creates the central term!

Cross terms $[\alpha_1 \cdot \alpha_1, [\cdot, \cdot]] = 0$

$[\cdot, \cdot], [\cdot, \cdot]$ create $4L_0$ (with $\alpha_{-1} \cdot \alpha_1$)

Any k : Exercise (Zwicklich p. 225, Baier-Lore p. 162)