Differential equations

- A differential equation (DE) describes how a quantity changes (as a function of time, position, ...).
 - A ball dropped from a building: $\frac{d}{dt}h(t) = -gt$

$$\frac{\mathrm{d}^2}{\mathrm{d}x^2}w(x) = \frac{S}{EI}w(x) + \frac{qx}{2EI}(x-L) \tag{2}$$

- Heat conduction:

. . .

$$-\frac{(h/2\pi)^2}{2m}\nabla^2\psi(\mathbf{r}) + V(\mathbf{r})\psi(\mathbf{r}) = E\psi(\mathbf{r})$$
(3)

$$\rho c_{P} \frac{\partial}{\partial t} T(\mathbf{r}, t) = \nabla \cdot (k \nabla T(\mathbf{r}, t))$$
(4)

- DEs can be classified in various ways:

Ordinary DEs (ODEs)

Only one independent variable Eqs (1) and (2)

Initial value problems(IVP)

Function value(s) known at certain values of the variable(s). Want function value(s) at arbitrary variable value(s). Eqs (1) and (4) In 1D: $f(t_0) = y_0$, compute f(t).

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Boundary value problems (BVP)

Many independent variables

Partial DEs (PDEs)

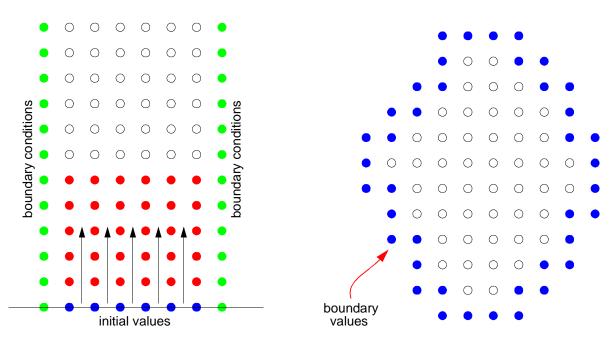
Eqs (3) and (4)

Function value(s) known on the domain borders. Want function values in other point(s) in the domain. Eqs (2), (3) and (4) In 1D: $f(x_1) = y_1$, $f(x_2) = y_2$, compute f(x).

(1)

Differential equations

• The difference between initial value problems and boundary value problems can be illustrated as below (according to NR, fig. 19.0.1).



- In general boundary value problems are harder to solve than initial value problems: BVPs need iteration IVP you 'only' need to integrate the equations in 'time' to reach the desired point.

- In this chapter we deal with ordinary differential equations both IVPs and BVPs.
- Probably there will be no time partial differential equations.

• The simplest ODE (IVP) is

$$\frac{dy}{dt} = f(t)$$

and its solution (starting from 'time' t_0)

$$y(t) = c + \int_{t_0}^{t} f(\tau) d\tau$$
. (Integration of the differential equation.)

- The solution y(t) is unique up to the constant c i.e. we get a family of solutions.
- The initial condition is needed to fix the value of the constant and to get a unique solution:

$$y(t_0) = y_0.$$

• The general form of an IVP is

$$\frac{dy}{dt} = f(t, y) \; .$$

Scientific computing III 2013: 13. Differential equations

Differential equations: initial value problems

• Solution methods for a single ODE are easily generalized to a group if ODEs (Use of vector notation makes it easier.)

$$\mathbf{y}(t) = \begin{bmatrix} y_1(t) \\ y_2(t) \\ \dots \\ y_n(t) \end{bmatrix}, \quad \mathbf{f}(t, \mathbf{y}) = \begin{bmatrix} f_1(t, \mathbf{y}) \\ f_2(t, \mathbf{y}) \\ \dots \\ f_n(t, \mathbf{y}) \end{bmatrix}, \quad \mathbf{y}_0 = \begin{bmatrix} A_1 \\ A_2 \\ \dots \\ A_n \end{bmatrix}$$

$$\mathbf{y}' = \mathbf{f}(t, \mathbf{y}), \ \mathbf{y}(t_0) = \mathbf{y}_0.$$

• ODEs having also higher derivatives can be transformed to groups of ODEs so that methods presented here can be applied to these higher order ODEs:

$$\begin{cases} y^{(n)} = f(t, y, y', \dots, y^{(n-1)}) \\ y(t_0) = A_1 \\ y'(t_0) = A_2 \\ \dots \\ y^{(n-1)}(t_0) = A_n \end{cases} \rightarrow \begin{cases} y_1' = y_2 \\ y_2' = y_3 \\ \dots \\ y_n' = f(t, y_1, y_2, \dots, y_n) \\ y_1(t_0) = A_1 \\ y_2(t_0) = A_2 \\ \dots \\ y_n(t_0) = A_n \end{cases}$$

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 $y' \equiv \frac{dy}{dx}$

dt

- · Numerical solution calls for discretization of the problem.
 - Instead of computing y(t) at any t we obtain values t_i in discrete points.
 (We talk about time here, but t can of course be any variable to be solved.)
 - The interval between points (in IVPs the time step) may vary as a function of time (variable time step).
- There several error sources in the integration of ODEs:
 - Discretization error: a too large time step \rightarrow decrease the time step.
 - Instability: small errors introduced in the beginning grow exponentially \rightarrow change the algorithm.
 - Stiff ODEs: phenomena with different time scales \rightarrow step size selected according to the fastest processes

Scientific computing III 2013: 13. Differential equations

Differential equations: initial value problems

- Something general can be said about the existence of the solution of the IVP y' = f(t, y)
 - A function f(t, y) is said to satisfy a Lipschitz condition in the variable y in the set $D \subset \mathbf{R}^2$ if $\exists L > 0$ so that $|f(t, y_1) f(t, y_2)| \leq L|y_1 y_2|$ whenever $(t, y_1), (t, y_2) \in D$.
 - L is called the Lipschitz constant of f.
 - A set $D \subset \mathbf{R}^2$ is said to be **convex** if whenever $(t_1, y_1), (t_2, y_2) \in D$ the point $((1 \lambda)t_1 + \lambda t_2, (1 \lambda)y_1 + \lambda y_2)$ also belongs to D for each λ , $0 \le \lambda \le 1$.
 - Suppose f(t, y) is defined in a convex set $D \subset \mathbf{R}^2$. If $\exists L > 0$, such that $\left| \frac{\partial}{\partial y} f(t, y) \right| \le L$, for all $(t, y) \in D$ then f satisfies the Lipschitz condition on D in variable y with Lipschitz constant L.

- Suppose that $D = \{(t, y) | a \le t \le b, -\infty < y < \infty\}$ and that f(t, y) is continuous on D. If f satisfies the Lipschitz condition on D in variable y then the IVP y' = f(t, y), $a \le t \le b$, $y(a) = \alpha$ has a unique solution y(t) for $a \le t \le b$.

- Example: $y' = 1 + t\sin(ty)$, $0 \le t \le 2$. y(0) = 0.

- Keeping *t* constant we apply the mean value theorem to function $f(y) = 1 + t\sin(ty)$:

$$\frac{df(\xi)}{dy} = \frac{f(y_2) - f(y_1)}{y_2 - y_1} , \quad y_1 < y_2, \quad y_1 < \xi < y_2.$$

$$\rightarrow \quad t^2 \cos(\xi t) = \frac{f(y_2) - f(y_1)}{y_2 - y_1} \rightarrow \quad 4 \ge \frac{|f(y_2) - f(y_1)|}{|y_2 - y_1|} \rightarrow \quad |f(y_2) - f(y_1)| \le 4|y_2 - y_1| .$$

- Thus L = 4 and the IVP has a unique solution.

- What about the stability of the IVP?
- The IVP y' = f(t, y), $a \le t \le b$, $y(a) = \alpha$ is said to be a well-posed problem (or a stable problem) if 1. A unique solution exists.
 - 2. For any $\varepsilon > 0$, $\exists k > 0$ with the property that, whenever $|\varepsilon_0| < \varepsilon$ and $|\delta(t)| < \varepsilon$,

a unique solution z(t) to the problem $z' = f(t, z) + \delta(t)$, $a \le t \le b$, $z(a) = \alpha + \varepsilon_0$ (PB) (=perturbed problem) exists with $|z(t) - y(t)| < k\varepsilon$, $\forall t, a \le t \le b$.

- One can show that the IVP y' = f(t, y) is a well-posed problem if

- 1) f(t, y) is a continuous function of t and y for all $(t, y) \in D$ and
- 2) it satisfies the Lipschitz condition $|f(t, y_1) f(t, y_2)| \le L|y_1 y_2|$, $(t, y_1), (t, y_2) \in D$.

Scientific computing III 2013: 13. Differential equations

Differential equations: initial value problems

- Example:

 $D = \{(t, y) | 0 \le t \le 1, -\infty < y < \infty\}$ y' = -y + t + 1 y(0) = 1.

- Since $\left|\frac{\partial(-y+t+1)}{\partial y}\right| = 1$, the function f(t, y) = -y+t+1 satisfies the Lipschitz condition.
- Morever, f is continuous \rightarrow the problem is well posed.
- Let's see what we get when we perturb the problem:

$$\frac{dz}{dt} = -z + t + 1 + \delta, \quad 0 \le t \le 1, \quad z(0) = 1 + \varepsilon_0$$

- Solutions to the original and perturbed problems are

$$y(t) = e^{-t} + t$$
, $z(t) = (1 + \varepsilon_0 - \delta)e^{-t} + t + \delta$.

- One can now easily see that if $|\delta| < \epsilon$ and $|\epsilon_0| < \epsilon$ then

$$|y(t) - z(t)| = \left| (\delta - \varepsilon_0) e^{-t} - \delta \right| \le |\varepsilon_0| + |\delta| |1 - e^{-t}| \le 2\varepsilon \quad \forall t.$$

- The above discussion concerns the stability of the problem itself.
- A problem is stable (well-posed) if the solution depends continuously on the initial value and on the right-hand side function *f*.
- The IVP could, however, be stable but ill-conditioned.
 - Condition number gives the sensitivity of the solution to the initial data.
 - A large condition number means an ill-conditioned problem.
- Assume we have a perturbed IVP: $y'(t;\varepsilon) = f(t, y(t;\varepsilon))$, $y(t_0;\varepsilon) = y_0 + \varepsilon$, $t \in [a, b]$.
 - Let's write $z(t;\varepsilon) = y(t;\varepsilon) y(t)$, where y(t) is the solution of the unperturbed IVP. This gives us

$$z'(t;\varepsilon) = f(t, y(t;\varepsilon)) - f(t, y(t)) \approx \frac{\partial f(t, y(t))}{\partial y} z(t;\varepsilon) \text{ and } z(t_0;\varepsilon) = \varepsilon.$$

- The above approximation is valid for small $\boldsymbol{\epsilon}$ (we have a well-posed problem).
- From the equation we can readily solve *z*:

$$z(t;\varepsilon) \approx \varepsilon \exp\left[\int_{t_0}^t \frac{\partial f(t', y(t'))}{\partial y} dt'\right]$$

- This means that if

$$\frac{\partial f(t, y(t))}{\partial y} \le 0 , \ t \in [a, b]$$

then $z(t;\varepsilon)$ probably remains bounded by ε as t increases. In this case the problem is well-conditioned.

- An example of an ill-conditioned problem is: $y' = \lambda y + g(t)$, $y(0) = y_0$, $\lambda > 0$.
- Now we can calculate $z(t;\varepsilon) = \varepsilon e^{\lambda t}$ indicating that the change in y(t) becomes increasing large as t increases.

Scientific computing III 2013: 13. Differential equations

Differential equations: initial value problems

- As another example of an ill-conditioned problem con-

sider

$$y' = 100y - 101e^{-t}, y(0) = 1.$$

- The solution is

$$y(t) = e^{-t}$$

- Perturbed problem:

$$y' \ = \ 100y - 101e^{-t}, \quad y(0) \ = \ 1 + \varepsilon \,, \label{eq:y}$$
 has the solution

$$y(t;\varepsilon) = e^{-t} + \varepsilon e^{100t}.$$

IVP:
$$y' = ay + be^{-t}$$

 $y(0) = y_0$
 $f(t, y) = ay + be^{-t}$
 $\frac{\partial f}{\partial y} = a$
Solution: $y(t) = \frac{1}{a+1}[(y_0 + ay_0 + b)e^{at} - be^{-t}]$

$$z(t;\varepsilon) \approx \varepsilon \exp\left[\int_{0}^{t} \frac{\partial f(t', y(t'))}{\partial y} dt'\right] = \varepsilon \exp\left[\int_{0}^{t} \frac{\partial (100y - 101e^{-t})}{\partial y} dt'\right] = \varepsilon e^{100t}$$

- Stability of the solution of an IVP y'(t) = f(t, y):
 - Solution is stable if $\forall \epsilon > 0 \ \exists \delta > 0$, such that every solution of the problem \hat{y} that fulfills

 $|y(0) - \hat{y}(0)| \le \delta$

also fulfills

 $|y(t) - \hat{y}(t)| \le \varepsilon$

for all $t \ge 0$.

- Solution is asymptotically stable if it is stable and

 $\lim_{t \to \infty} |y(t) - \hat{y}(t)| = 0.$

- A simple example: $y' = \lambda y$. Solution is $y(t) = y(0)e^{\lambda t}$.
 - If $\lambda > 0$ all solutions grow and diverge from each other exponentially: unstable solution.
 - If $\lambda < 0\,$ all solution converge to zero: asymptotically stable solution.
 - If λ is complex ($\lambda = a + ib$) then stability depends on Re(λ): $y(t) = y(0)e^{(a+ib)t}$.
 - If $Re(\lambda) = 0$ then we get an oscillatory solution that is stable but no asymptotically stable.
- In addition to the stability definitions given above, the numerical method used to solve the IVP can also be stable or unstable.

Scientific computing III 2013: 13. Differential equations

Differential equations: initial value problems

- Euler's method is the simplest possible method to solve an IVP.
- It is not used much in practical problems but its derivation illustrates the concepts needed in deriving more advanced methods.
- The IVP we are trying to solve is the same as before

$$\frac{dy}{dt} = f(t, y), \quad a \le t \le b, \ y(a) = \alpha.$$

• The problem is discretized by computing the y values in mesh points $\{t_0, t_1, ..., t_N\}$ in interval [a, b]:

$$t_i = a + ih, \quad i = 0, 1, \dots, N.$$

- The distance between the consecutive points h = (b-a)/N is called the step size.
- Assume the solution y(t) has two continuous derivatives.
- We will expand it as a Taylor series

$$\begin{split} y(t_{i+1}) &= y(t_i) + (t_{i+1} - t_i)y'(t_i) + \frac{1}{2}(t_{i+1} - t_i)^2 y''(\xi_i), \quad \xi_i \in [t_i, t_{i+1}] \\ \text{or} \quad y(t_{i+1}) &= y(t_i) + hy'(t_i) + \frac{1}{2}h^2 y''(\xi_i). \end{split}$$

- Since y(t) is the solution to the IVP we can write the expansion as

$$y(t_{i+1}) = y(t_i) + hf(t_i, y(t_i)) + \frac{1}{2}h^2y''(\xi_i)$$
.

- In Euler's method we simply drop the 2nd order term and get the iteration formula

$$y_0 = \alpha$$
,
 $y_{i+1} = y_i + hf(t_i, y_i)$, $i = 0, 1, ..., N$.

- So we get an approximation to the solution

 $y_i \approx y(t_i)$.

- The algorithm is extremely simple

1. Set h = (b-a)/N, t = a, $y = \alpha$. 2. For i = 1, 2, ..., N do Set $y \leftarrow y + hf(t, y)$. Set $t \leftarrow a + ih$.

Scientific computing III 2013: 13. Differential equations

Differential equations: initial value problems

- As a simple example take

$$\frac{dy(t)}{dt} = \sin t \, \cdot \, y(0) = 0 \, \cdot \, 0 \le t \le 3 \, .$$
- The exact solution is

$$y(t) = 1 - \cos t.$$

- The Euler algorithm in awk:

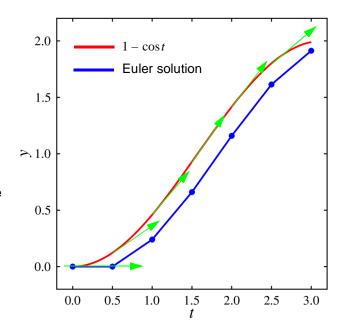
```
BEGIN {
    h=0.5; y=0; tm=3.0; n=tm/h;
    print 0,y;
    for (i=0;i<n;i++) {
        t=i*h;
        y=y+h*sin(t);
        t=(i+1)*h
        print t,y;
    }
}</pre>
```

- In the figure one can see the graphical interpretation of the alrgorithm.
 - When the problem is well-posed

 $f(t_i, y_i) \approx y'(t_i) = f(t_i, y(t_i))$.

- Euler's method is easily generalized to groups of equations:

$$\mathbf{y}_{k+1} = \mathbf{y}_k + h\mathbf{f}(t_k, \mathbf{y}_k).$$



- What about the error behavior of the Euler's method?
 - By subtracting the exact solution from the Euler iteration result we get

$$y(t_{k+1}) - y_{k+1} = y(t_k) - y_k + h[f(t_k, y(t_k)) - f(t_k, y_k)] + \frac{h^2 y''(\xi_k)}{2} , \quad \xi_k \in [t_k, t_{k+1}].$$
(1)

- The left hand side of this equation is the total error of the approximation y_{k+1} and it consists of
 - 1. Errors cumulated from the previous step at t_k
 - 2. Error formed in computing the current step (local error).
- Error 2 is called local because if we happened to get the exact value in the previous step (k) then the only error that is left is $h^2 y''(\xi_k)/2 \equiv l_{k+1}$.
- In general the local error is defined as

$$l_n = \hat{y}(t_n) - y_n$$

and \hat{y} is the solution of the local problem:

$$\hat{y}'(t) = f(t, \hat{y}(t))$$

$$\hat{y}(t_{n-1}) = y_{n-1} \quad .$$

- Due to the local error $h^2 y''(\xi_k)/2 = O(h^2)$ Euler's method is of the order one.

- In general the order of a method is p if the local error behaves as $O(h^{p+1})$.

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Differential equations: initial value problems

- Using the mean value theorem we can estimate the term in the brackets in (1):

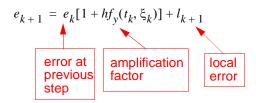
$$h[f(t_k, y(t_k)) - f(t_k, y_k)] = h \frac{\partial f(t_k, \xi_k)}{\partial y} [y(t_k) - y_k] \quad , \quad \xi_k \in [y(t_k), y_k].$$

- Substituting this to eq. (1) we get

$$\begin{split} y(t_{k+1}) - y_{k+1} &= y(t_k) - y_k + h \frac{\partial f(t_k, \xi_k)}{\partial y} [y(t_k) - y_k] + \frac{h^2 y''(\xi_k)}{2} \\ &= [y(t_k) - y_k] [1 + h f_y(t_k, \xi_k)] + \frac{h^2 y''(\xi_k)}{2} \end{split}$$

where $f_y(t_k, \xi_k) &= \frac{\partial f(t_k, \xi_k)}{\partial y}$

- So we get the total error at step k + 1 as



- As far as $|1 + hf_y(t_k, \xi_k)| < 1$ the total error in Euler's method does not grow. - If $|1 + hf_y(t_k, \xi_k)| > 1$ the error grows during the iteration without limits.

- Let's look at the test problem

 $y' = \lambda y$,

- where $\boldsymbol{\lambda}$ is a complex constant.
- The solution is $y(t) = y(0)e^{\lambda t}$.
- Applying Euler's method we get

 $y_k = y_{k-1} + h\lambda y_{k-1} = (1 + h\lambda)y_{k-1} = \dots = y(0)(1 + h\lambda)^k.$

- Let now $\text{Re}(\lambda) < 0$ so that the solutions are asymptotically stable and approach zero.
- This means that the numerical approximations $|y_k|$ must not grow.

In Euler's method this means $|1 + h\lambda| \le 1$.

- The absolute stability region D of a numerical method consists of those
 - complex numbers $z = h\lambda$ that fulfill the condition

$$|y_k| \le |y_{k-1}|, \quad k = 1, 2, \dots$$

- Absolute stability interval is the part of the real axis that belongs to D.
- The absolute stability region of Euler's method is the circle $D = \{z \in \mathbb{C} | |z+1| \le 1\}.$ and the interval is [-2, 0].

Scientific computing III 2013: 13. Differential equations

Generally for the Euler's method: $\left|1+hf_y\right|<1~.$

Differential equations: initial value problems

- Example: solving
 - y' = -100y + 100, $y(0) = y_0$.

using Euler's method.

- Exact solution is $y(t) = (y_0 1)e^{-100t} + 1$ and it is clearly stable.
- Euler's method gives a difference equation

 $y_{k+1} = y_k + h(-100y_k + 100) = (1 - 100h)y_k + 100h$

with solution

$$y_k = (y_0 - 1)(1 - 100h)^k + 1.$$

- Let the initial value be $y_0 = 2$.
- Exact solution is now $y(t) = e^{-100t} + 1$.
- Euler's method gives $y_k = (1 100h)^k + 1$
- We see immediately that if h > 0.02 the solution $|y_k|$ grows without limit when k increases.

- The error formula of Euler's method contains *h* both in the error amplification term and in the local error term:

$$e_{k+1} \,=\, e_k [1 + h f_y(t_k, \xi_k)] + \frac{h^2 y''(\xi_k)}{2} \;. \label{eq:ek+1}$$

- In practical computations h should be

as large as possible in order to speed up the calculations. so small that the error amplification factor < 1 so small that the local error term is acceptable.

- The derivatives are difficult to compute but the local term can be estimated based on the iterations:

$$y_{a}'' = \frac{y_{k}' - y_{k-1}'}{t_{k} - t_{k-1}'}.$$

- By requiring that the local error fulfills the condition

$$\frac{|y_a''|h^2}{2} < \varepsilon$$

we get the upper limit to the step size

$$h < \sqrt{\frac{2\varepsilon}{|y_a''|}}$$
.

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Differential equations: initial value problems

- Even if it would be possible from the point of view of the consumed CPU time to decrease the step size indefinitely we wouldn't get better and better results: finally the round-off errors show up.
- The figure on the right shows the relative error of the solution of the IVP

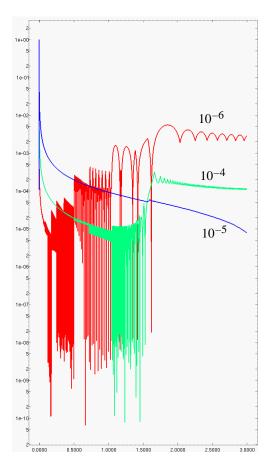
 $y' = \sin t$. y(0) = 0. $0 \le t \le 3$

using step sizes

 $h = 10^{-4}, 10^{-5}, 10^{-6}$

and using single precision floating point numers.

- You probably remember the same phenomenon in computing the derivative of a function using simple difference quotient.



• We can go further in the Taylor's series

$$y(t_{i+1}) = y(t_i) + (t_{i+1} - t_i)y'(t_i) + \frac{1}{2}(t_{i+1} - t_i)^2y''(t_i) + \dots$$

and develop higher order methods.

- These methods have the advantage of higher-order local truncation error.
- However, one has to calculate the higher derivatives $f^{(n)}(t, y)$ which is time consuming in many practical problems.
- An alternative to the Euler's method is an implicit version of it (backward Euler's method):
 - Integrate y'(t) = f(t, y) over one time step

$$y(t_{k+1}) = y(t_k) + \int_{t_k}^{t_{k+1}} f(t, y(t)) dt$$

- By approximating the intergral by $(t_{k+1} t_k)f(t_k, y_k)$ we get the Euler's method.
- However, we can also use $(t_{k+1} t_k)f(t_{k+1}, y_{k+1})$ when we get the backward Euler:

$$y_{k+1} = y_k + hf(t_{k+1}, y_{k+1})$$
.

- The word implicit means that new point y_{k+1} is on both sides of the iteration formula. It has to be solved by numerical methods.
- One can show that this methods has a error amplification factor $(1 hf_y)^{-1}$ which means that when applied to a stable linear problem with constant coefficients when $\text{Re}(\lambda) \le 0$, we get $|1 h\lambda|^{-1} \le 1$ for all h > 0.

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Differential equations: initial value problems

- The methods described above were presented to illustrate the important concepts in numerical integration of IVPs.
- They are seldom used in practical applications.
- Probably the most often used method to solve IVPs is the Runge-Kutta (RK) method (or one of the many RK methods).
- It offers the higher orders without the need to compute derivatives of f.
- The only drawback is that the number of function evaluations of *f* per time step may be large.
- We need a couple of results concerning the Taylor's theorem of a function of two variables.
 - Suppose f(y, t) has all its partial derivatives up to n + 1 continuous on $D = \{(t, y) | a \le t \le b, c \le y \le d\}$.

$$\cdot$$
 Let $(t_0, y_0) \in D$.

- For every $(t, y) \in D \ \exists (\xi \in [t, t_0], \eta \in [y, y_0])$ with

$$f(t, y) = P_n(t, y) + R_n(t, y)$$

$$\begin{split} P_n(t,y) &= f(t_0,y_0) + \left[(t-t_0) \frac{\partial f(t_0,y_0)}{\partial t} + (y-y_0) \frac{\partial f(t_0,y_0)}{\partial y} \right] \\ &+ \left[\frac{(t-t_0)^2}{2} \frac{\partial^2}{\partial t^2} f(t_0,y_0) + \frac{(y-y_0)^2}{2} \frac{\partial^2}{\partial y^2} f(t_0,y_0) + (t-t_0)(y-y_0) \frac{\partial^2}{\partial t \partial y} f(t_0,y_0) \right] \\ &+ \dots + \left[\frac{1}{n!} \sum_{j=0}^n \binom{n}{j} (t-t_0)^{n-j} (y-y_0)^j \frac{\partial^n}{\partial t^{n-j} \partial y^j} f(t_0,y_0) \right] \end{split}$$

$$R_{n}(t,y) = \frac{1}{(n+1)!} \sum_{j=1}^{n+1} {n+1 \choose j} (t-t_{0})^{n+1} (y-y_{0})^{j} \frac{\partial^{n+1}}{\partial t^{n+1-j} \partial y^{j}} f(\xi,\eta).$$

- In other words P_n is the Taylor polynomial and R_n is the remainder term.

- Let's see how by using the RK method we could increase the order of the Euler's method by one.
- This means that we have to approximate the term

$$T^{(2)}(t, y) = f(t, y) + \frac{h}{2}f(t, y)$$

with error no greater than $O(h^2)$.

- We take the approximation to have the form $a_1 f(t + \alpha_1, y + \beta_1)$.

- Since
$$f(t, y) = \frac{df(t, y)}{dt} = \frac{\partial}{\partial t}f(t, y) + \frac{\partial}{\partial y}f(t, y)y'(t)$$
 and $y'(t) = f(t, y)$
 $\rightarrow T^{(2)}(t, y) = f(t, y) + \frac{h}{2}\frac{\partial}{\partial t}f(t, y) + \frac{h}{2}\frac{\partial}{\partial y}f(t, y)y'(t)$.

- Now we expand $f(t + \alpha_1, y + \beta_1)$ into a Taylor polynomial of degree one about (t, y)

$$\begin{aligned} a_1 f(t+\alpha_1, y+\beta_1) &= a_1 f(t, y) + a_1 \alpha_1 \frac{\partial}{\partial t} f(t, y) + a_1 \beta_1 \frac{\partial}{\partial y} f(t, y) + a_1 R_1 (t+\alpha_1, y+\beta_1) \\ R_1 (t+\alpha_1, y+\beta_1) &= \frac{\alpha_1^2}{2} \frac{\partial^2}{\partial t^2} f(\xi, \eta) + \frac{\beta_1^2}{2} \frac{\partial^2}{\partial y^2} f(\xi, \eta) + \alpha_1 \beta_1 \frac{\partial^2}{\partial t \partial y} f(\xi, \eta) \,. \end{aligned}$$

- The coefficients a_1 , α_1 , and β_1 can be determined by setting

$$f(t, y) + \frac{h}{2}f(t, y) = a_1f(t + \alpha_1, y + \beta_1)$$

$$\Rightarrow \qquad f(t, y) + \frac{h}{2\partial t}f(t, y) + \frac{h}{2\partial y}f(t, y)y'(t) = a_1f(t, y) + a_1\alpha_1\frac{\partial}{\partial t}f(t, y) + a_1\beta_1\frac{\partial}{\partial y}f(t, y) \quad \text{(remember } y' = f(t)$$

and matching the coefficients of f and its derivatives.

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Differential equations: initial value problems

- The results is

$$a_1 = 1, \ \alpha_1 = \frac{h}{2}, \ \beta_1 = \frac{h}{2}f(t, y).$$

and

$$T^{(2)}(t, y) = f\left(t + \frac{h}{2}, y + \frac{h}{2}f(t, y)\right) - R_1\left(t + \frac{h}{2}, y + \frac{h}{2}f(t, y)\right) \quad .$$

- If the 2nd partial derivatives of f are bounded then the remainder term R_1 is of order $O(h^2)$.

- The method derived above is a 2nd order (and two-stage) RK method and the algorithm itself is simple:

1. $y_0 = y(0)$.

2.
$$y_{i+1} = y_i + hf\left(t_i + \frac{h}{2}, y_i + \frac{h}{2}f(t_i, y_i)\right)$$
.

- This has a name of its own: the midpoint method.

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$$y_0 = y(t_0)$$

$$y_{i+1} = y_i + hT^{(n)}(t_i, y_i)$$

$$T^{(n)} = f(t_i, y_i) + \frac{h}{2}f(t_i, y_i)$$

$$+ \dots + \frac{h^{n-1}}{(n+1)!}f^{(n-1)}(t_i, y_i)$$

)

- However, the most used RK method is the 4th order version:

$$\begin{split} y_0 &= y(0) \\ k_1 &= hf(t_i, y_i) \\ k_2 &= hf\left(t_i + \frac{h}{2}, y_i + \frac{k_1}{2}\right) \\ k_3 &= hf\left(t_i + \frac{h}{2}, y_i + \frac{k_2}{2}\right) \\ k_4 &= hf(t_{i+1}, y_i + k_3) \\ y_{i+1} &= y_i + \frac{1}{6}(k_1 + 2k_2 + 2k_3 + k_4). \end{split}$$

- Note that in this algorithm the function *f* is evaluated 4 times per time step.
- In general a s-stage RK method can be written as

$$y_{i+1} = y_i + h \sum_{j=1}^{s} b_j k_j$$
$$k_j = f\left(t_i + c_j h, y_i + h \sum_{l=1}^{2} a_{jl} k_l\right)$$
$$c_j = \sum_{l=1}^{s} a_{jl}.$$

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Differential equations: initial value problems

- One can describe the RK method as follows:
 - 1. Take an Euler step of length c_2h from (t_i, y_i) . The derivative at the new point is k_2 .
 - 2. Return to the original point and take a new Euler step (c_3h) using the weighted average
 - of k_1 and k_2 as the derivative. We get a new derivative k_3 .
 - In this way a number of approximation are computed for the derivative and they all are used in the end to compute the step taken from (t_i, y_i) .
 - The number of stages in a RK method is not directly the order of the method:

Order	1	2	3	4	5	6	7	8	9	10
Stages	1	2	3	4	6	7	9	11	$12 \le s \le 17$	$13 \le s \le 17$

- The table above explains the popularity of the four-stage RK method.

- Most practical routines for integration of an IVP require the user to give an upper limit to the error in the end result.
 - The program then adjusts the time step h accordingly.
 - This means that the algorithm somehow has to be able to estimate the error.
 - One way to do this is to compute the solution at the same time with step sizes h and 2h.
 - Assuming that the local error introduced by the larger step is twice that of the smaller step one can get an estimate for the error by

$$\left|l_{n}\right| \approx \frac{1}{2} \frac{\left|\tilde{y}_{n} - y_{n}\right|}{2^{p} - 1},$$

where \tilde{y}_n and y_n are the solution at time *n*.

- However, computing with two time steps is an extra effort.
- In the RK method the error estimate can be computed by solving the IVP with two RK solvers simultaneously.
- The two solvers can be chosen in such a way that the derivatives needed by the lower order solver are already computed by the higher order method.
- This pair of solvers is called the embedded Runge-Kutta pair.
- They are sometimes also called Runge-Kutta-Fehlberg methods.
- If the RK method is written in the form $y_{n+1} = y_n + F(t_n, y_n, h; f)$ then the method is stable if F fulfills the Lipschitz condition

$$|F(t, y, h;f) - F(t, z, h;f)| \le L|y - z|$$
,
for all $a \le t \le b$, $-\infty < y, z < \infty$.

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Differential equations: initial value problems

- Below is a comparison of the maximum error of various one-step methods¹.
 - A one-step method is a method that only uses y_i to compute the next step y_{i+1} .

- The IVP was
$$y' = -5ty^2 + \frac{5}{t} - \frac{1}{t^2}$$
, $y_0 = 1$, $t \in [1, 10]$

- The exact solution is y(t) = 1/t.

h	Euler	Impl. Euler	Trapetsi	RK2	RK4
0.2	3.33e-2	9.23e-3	1.39e-3	1.39e-3	5.34e-3
0.1	9.09e-3	5.21e-3	2.83e-4	2.83e-4	2.04e-4
0.05	3.43e-3	2.80e-3	7.01e-5	7.01e-5	9.65e-6
0.02	1.27e-3	1.17e-3	1.11e-5	1.11e-5	2.09e-7
0.01	6.20e-4	5.97e-4	2.77e-6	2.77e-6	1.23e-8
0.005	3.08e-4	3.01e-4	6.92e-7	6.92e-7	7.47e-10
0.002	1.22e-4	1.21e-4	1.11e-7	1.11e-7	1.88e-11

- The trapezoidal method is $y_{k+1} = y_k + \frac{h}{2}[f(t_{k+1}, y_{k+1}) + f(t_k, y_k)]$.

^{1.} From J. Haataja et al., Numeeriset menetelmät käytännössä, CSC, 1999

- In multistep (MS) methods one uses also information from earlier time steps than only the previous one.
 - I.e. we want to use also

$$y_{k-1}, y_{k-2}, \dots$$

 $f(t_{k-1}, y_{k-1}), f(t_{k-2}, y_{k-2}), \dots$

to approximate the next point y_{k+1} .

- In general a *r* step linear multistep method can be written as $f_i \equiv f(t_i, y_i)$

$$\sum_{j=-r+1}^{1} \alpha_{j} y_{k+j} = h \sum_{j=-r+1}^{1} \beta_{j} f_{k+j}$$

where the coefficients fulfill the conditions

$$\alpha_1 = 1, \, \alpha_{-r+1} \neq 0, \, \beta_{-r+1} \neq 0.$$

- Methods can be derived by approximating the integral in

$$y(t_{k+1}) = y(t_k) + \int_{t_k}^{t_{k+1}} f(t, y(t)) dt$$

with a polynomial

$$P(t) = a_0 + a_1 t + a_2 t^2 + \ldots + a_s t^s$$
 and by integrating it.

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Differential equations: initial value problems

- The unknown coefficients a_i are determined by fixing the polynomial to certain values y_i and f_i , i = k + 1, k, k 1, ...
- The good point in the MS methods is that the function f is only evaluted once per time step.
- On the other hand, in the beginning of the iteration there are no previous data points available. - These can be generated by perfoming a couple of steps using a one-step method.
- Also, changing the time step is not so straightforward in MS methods.
- MS methods are divide to

explicit (open) methods, $\beta_1 = 0$ and

implicit (closed) methods, $\beta_1 \neq 0$.

- Of the explicit MS methods the Adams-Bashford (AB) is maybe the most well known.
 - It is of the form

$$y_{k+1} = y_k + h \sum_{j=-r+1}^{0} \beta_j f_{k+j}.$$

- The unknown new iteration point does not show up in the RHS of the equation so this is an open method.
- Below are the five first AB methods

$$\begin{split} y_{k+1} &= y_k + hf_k \\ y_{k+1} &= \frac{h}{2}[3f_k - f_{k-1}] \\ y_{k+1} &= y_k + \frac{h}{12}[23f_k - 16f_{k-1} + 5f_{k-2}] \\ y_{k+1} &= y_k + \frac{h}{24}[55f_k - 59f_{k-1} + 37f_{k-2} - 9f_{k-3}] \\ y_{k+1} &= y_k + \frac{h}{720}[1901f_k - 2774f_{k-1} + 2616f_{k-2} - 1274f_{k-3} + 251f_{k-4}] \end{split}$$

- One can show that an AB method with r steps is of the order $O(h^r)$.

- The stability interval of the AB methods are rather narrow, so that they are not suitable for stiff problems.

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Differential equations: initial value problems

• The implicit Adams-Moulton (AM) method is more suitable for stiff problems:

$$y_{k+1} = y_k + h \sum_{j=-r+2}^{1} \beta_j f_{k+j}.$$

- The five first AM methods are

$$\begin{aligned} y_{k+1} &= y_k + hf_{k+1} \\ y_{k+1} &= y_k + \frac{h}{2}[f_{k+1} + f_k] \\ y_{k+1} &= y_k + \frac{h}{12}[5f_{k+1} + 8f_k - f_{k-1}] \\ y_{k+1} &= y_k + \frac{h}{24}[9f_{k+1} + 19f_k - 5f_{k-1} + f_{k-2}] \\ y_{k+1} &= y_k + \frac{h}{720}[251f_{k+1} + 646f_k - 264f_{k-1} + 106f_{k-2} + 19f_{k-3}] \end{aligned}$$

- When $r \ge 2$ the order of the AM method is r + 1.
- The absolute stability intervals of the methods are
 - 1 $[-\infty, 0]$
 - 2 [-6, 0]
 - 3 [-8, 0]
 - 4 [-1.8, 0]

- The AM methods give better results than the AB methods.
 - However, the drawback of the AM method is that due to its implicity one has to solve the new iteration point from a group of equations.
- A combination method often used is so called predictor-corrector method.
 - In it the AB method produces a initial prediction to the new y value while the AM method is used to correct it more precise.
 - The correction may be done many times but not until the final convergence.
 - For example the pair of 3rd degree:

r

$$\begin{cases} y_{k+1}^{p} = y_{k} + \frac{h}{12} [23f_{k} - 16f_{k-1} + 5f_{k-2}] \\ y_{k+1}^{c} = y_{k} + \frac{h}{12} [5f(t_{k+1}, y_{k+1}^{p}) + 8f_{k} - f_{k-1}] \end{cases}$$

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Differential equations: initial value problems

- So called backward differentiation formulae (BDF) methods are well suited for stiff ODEs.
 - They are derived by requiring a polynomial going through the previous iteration points:

$$P_r(t) = \sum_{j=-1}^{r-1} y(t_{n-j}) l_{jn}(t)$$
, where $l_{jn}(t)$ is the Lagrange interpolation polynomial.

- Because P_r is (an approximate) solution

$$P'_{r}(t_{n+1}) = y'(t_{n+1}) = f(t_{n+1}, y(t_{n+1}))$$

- Combining the above equations and solving $y(t_{n+1})$ we get

$$y_{n+1} = \sum_{j=0}^{r-1} \alpha_j y_{n-j} + h\beta f(t_{n+1}, y_{n+1}),$$

where α_i and β are obtained by expanding the Lagrange interpolation polynomial in the first equation above.

- The method is implicit and the order is r.
- Below are the five first BDF methods:

$$y_{k+1} = y_k + hf_{k+1}$$
 (This is backward Euler.)

$$y_{k+1} = \frac{4}{3}y_k - \frac{1}{3}y_{k-1} + \frac{2}{3}hf_{k+1}$$

$$y_{k+1} = \frac{18}{11}y_k - \frac{9}{11}y_{k-1} + \frac{2}{11}y_{k-2} + \frac{6}{11}hf_{k+1}$$

$$y_{k+1} = \frac{48}{25}y_k - \frac{36}{25}y_{k-1} + \frac{16}{25}y_{k-2} - \frac{3}{25}y_{k-3} + \frac{12}{25}hf_{k+1}$$

$$y_{k+1} = \frac{300}{137}y_k - \frac{300}{137}y_{k-1} + \frac{200}{137}y_{k-2} - \frac{75}{137}y_{k-3} + \frac{12}{137}y_{k-4} + \frac{60}{137}hf_{k+1}$$

- Stiff ODEs are those where there are two or more disparate time scales.
 - An example: y' = -100y + 100t + 101, $y(0) = y_0$.

- Solution:
$$y(t) = (y_0 - 1)e^{-100t} + t + 1$$

- By varying the initial value y_0 slightly causes large deviations in the initial behavior of the solution.
- Euler method is results very sensitive to the initial value: even a small error during the iteration causes the iteration to pick a solution with a different y_0 ; par-

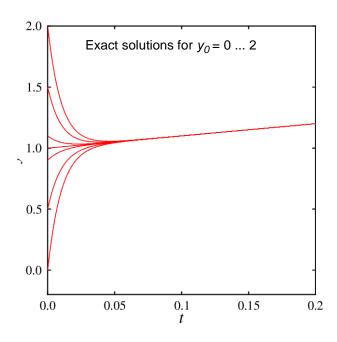
ticularly one with $y_0 \neq 1$.

- Morever, with large enough time steps we get outside the stability interval.
- Demo.
- Stiffness is difficult to define exactly:

1) Process described by the ODE contains disparate timescales.

2) A group of ODEs is stiff if the eigenvalues of the Jacobian **J** ($J_{ij} = \partial f_i / \partial y_j$) has eigenvalues that differ greatly in magnitude.

3) When using an explicit method a much smaller time step must be used as which would be dictated by the accuracy criteria.



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Differential equations: initial value problems

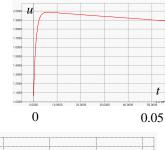
- Example (from NR, section 16.6)

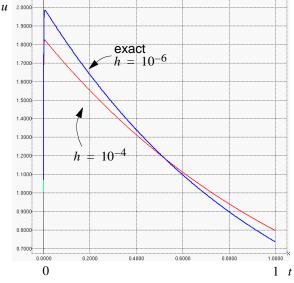
$$\begin{cases} u' = 998u + 1998v \\ v' = -999u - 1999v \end{cases}$$

- Exact solution

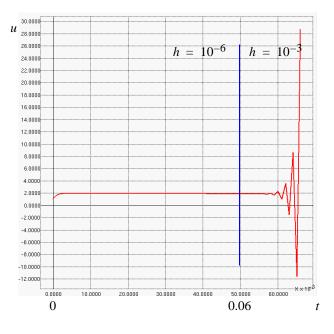
 $\begin{cases} u = 2e^{-t} - e^{-1000t} \\ v = -e^{-t} + e^{-1000t} \end{cases}$

- Using Euler's method we have to use really small steps:
- Although the smoothness of the function would allow a longer time step outside the initial transient behavior (term e^{-1000t})
- What about having a small step in the beginning and increasing it at say $t > 10^{-2}$





- Below is the results of a version of Euler which changes the stepsize after a certain time:



- This does not help because the instability of the transient is still there though the solution itself is a smoothly varying function.
- Or to put it in another way: In order the solution to be stable we must use a smaller time step than the one based only on the accuracy criteria.

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Differential equations: initial value problems

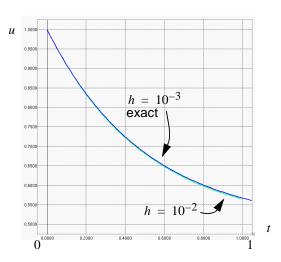
- For comparison consider the following equation that is not stiff but has a rather similar solution:

$$\begin{cases} u' = -u + v \\ v' = u - v \end{cases}$$

- And the solution with initial conditions u(0) = 1, v(0) = 0 is

$$\begin{cases} u = \frac{1}{2}(1 + e^{-2t}) \\ v = \frac{1}{2}(1 - e^{-2t}) \end{cases}$$

- Using Euler with $h = 10^{-3}, 10^{-2}$ gives reasonable results:



- For stiff ODEs the implicit methods work better.
- Take a simple example:

$$y' = -cy, c > 0.$$

- Euler's iteration gives us

$$y_{n+1} = y_n + hy'_n = (1 - hc)y_n$$

- The method is unstable if $h > \frac{2}{c}$.
- Now using implicit method (backward Euler) we get

$$y_{n+1} = y_n + hy'_{n+1}$$
 or
 $y_{n+1} = \frac{y_n}{1+ch}.$

- This method is stable for all values of h.
- So with implicit methods we can use larger time steps in considering the stability of the method.
- However, the accuracy may require to use smaller stepsize.

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Scientific computing III 2013: 13. Differential equations
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Differential equations: initial value problems

• Practical tools:

GSL: routines for many stepping algorithms and for timestep control (including various Runge-Kutta and implicit methods)
Matlab: also many routines: say help funfun to get a list.
Maxima: symbolic solver, example solving the last example equation:

```
Maxima 5.9.1 http://maxima.sourceforge.net

Using Lisp CMU Common Lisp 19a

Distributed under the GNU Public License. See the file COPYING.

Dedicated to the memory of William Schelter.

This is a development version of Maxima. The function bug_report()

provides bug reporting information.

(%i1) e1:'diff(u(t),t)=-u(t)+v(t);

(%o7) \frac{d}{dt}u(t) = v(t) - u(t)

(%i8) e2:'diff(v(t),t)=u(t)-v(t);

(%o8) \frac{d}{dt}v(t) = u(t) - v(t)

(%i9) desolve([e1,e2],[u(t),v(t)]);

(%o9) \left[u(t) = \frac{v(0) + u(0)}{2} - \frac{(v(0) - u(0))e^{-2t}}{2}, v(t) = \frac{(v(0) - u(0))e^{-2t}}{2} + \frac{v(0) + u(0)}{2}\right]

(%i10) |
```

- Practical tools (continued)
 - SLATEC library includes (at least) the following algorithms:
 - DDEABM: Adams-Bashforth method
 - DDEBDF: backward differentiation; for stiff problems
 - DDERKF: Runge-Kutta-Fehlberg method
 - User should, in addition, provide the function that calculates the RHS of the equation(s) and for the backward differentiation method Jacobian J of the problem needed:

$$\begin{cases} y_1' = f_1(t, y_1, \dots, y_N) \\ y_2' = f_2(t, y_1, \dots, y_N) \\ \dots \\ y_N' = f_N(t, y_1, \dots, y_N) \end{cases} \qquad \mathbf{J} = \begin{bmatrix} \frac{\partial f_1}{\partial y_1} \frac{\partial f_1}{\partial y_2} \cdots \frac{\partial f_1}{\partial y_N} \\ \frac{\partial f_2}{\partial y_1} \frac{\partial f_2}{\partial y_2} \cdots \frac{\partial f_2}{\partial y_N} \\ \dots \\ \frac{\partial f_N}{\partial y_1} \frac{\partial f_N}{\partial y_2} \cdots \frac{\partial f_N}{\partial y_N} \end{bmatrix}$$

• What method to use?

For non-stiff ODEs explicit RK, or predictor-corrector and AM if computing the derivative is expensive. For stiff ODEs BDF and implicit RK.

- Now a nice demo.

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Differential equations: boundary value problems¹

• Boundary value problems (BVPs) are harder to solve because the conditions have to be fulfilled in more than one point. - The general form of 1D BVP is

 $\begin{cases} \mathbf{y}' = \mathbf{f}(x, \mathbf{y}) \\ \mathbf{g}(\mathbf{y}(a), \mathbf{y}(b)) = 0 \end{cases}$ (y and g have N components)

0

- Linear boundary conditions (BC) can be written in the form $B_a \mathbf{y}(a) + B_b \mathbf{y}(b) = \mathbf{b}$.

- If the BCs are separable they can be written as

$$\mathbf{y}(a) = \mathbf{b}_a$$
$$\mathbf{y}(b) = \mathbf{b}_b.$$

- Sometimes the DE and BC contain a parameter λ :

$$\begin{cases} \mathbf{y}' = \mathbf{f}(x, \mathbf{y}, \lambda) \\ \mathbf{g}(\mathbf{y}(a), \mathbf{y}(b), \lambda) = \end{cases}$$

and the objective is to determine those values of λ for which the system has non-trivial solutions (eigenvalues and eigenfunctions).

- E.g. Schrödinger equation

$$\left[-\frac{1}{2}\nabla^2 + V(\mathbf{r})\right]\phi(\mathbf{r}) = E\phi(\mathbf{r})$$

- In 1D a particle constrained by infinite potential walls:

$$\begin{cases} -\frac{1}{2}\phi''(x) + V(x)\phi(x) = E\phi(x) \\ \phi(a) = 0 \\ \phi(b) = 0 \end{cases}$$
, which can be written as a system of 1st order DEs
$$\begin{cases} \psi(x) = 2[V(x) - E]\phi(x) \\ \phi(x) = \psi'(x) \end{cases}$$

^{1.} Sources: J. Haataja et al., Numeeriset menetelmät käytännössä, CSC, 1999; R. L. Burden, J. D. Faires, Numerical Analysis, PWS-KENT, 1989.

- In the shooting method (SM) we utilize methods developed for IVPs.
 - We integrate the DE starting from a chosen initial condition:

$$\mathbf{y}' = \mathbf{f}(x, \mathbf{y}), \ \mathbf{y}(a) = \mathbf{s}.$$

- We have to choose s so that the solution $\mathbf{y}(x;\mathbf{s})$ fulfills the BC $\mathbf{g}(\mathbf{y}(a;\mathbf{s}), \mathbf{y}(b;\mathbf{s})) = 0$.
- The solution of the original BVP is thus $\mathbf{y}(x;\mathbf{s}^*)$ where \mathbf{s}^* is the solution of the group of equations

 $\mathbf{F}(\mathbf{s}) = \mathbf{g}(\mathbf{s}, \mathbf{y}(b; \mathbf{s})) = 0.$

- This generally a non-linear group of equations.

Scientific computing III 2013: 13. Differential equations

Differential equations: boundary value problems

- A common form of BVPs is the 2nd order equation

$$y'' = f(x, y, y'), a \le x \le b,$$

$$y(a) = \alpha,$$

$$y(b) = \beta.$$

- This is converted to a pair of 1st order DEs as $(y \rightarrow y_1, y' \rightarrow y_2)$

$$\begin{cases} y_1' = y_2 \\ y_2' = f(x, y_1, y_2) \end{cases}$$

- Existence of the solution:

- Assume that the function f and its derivatives $\partial f/\partial y$, $\partial f/\partial y'$ are continuous in
 - $D = \{(x, y, y') | a \le x \le b, -\infty \le y \le \infty, -\infty \le y' \le \infty\}.$
- Now, if for all for all $(x, y, y') \in D$ the following apply

a)
$$\frac{\partial}{\partial y} f(x, y, y') > 0$$
 and

b)
$$\exists M$$
 such that $\left| \frac{\partial}{\partial y'} f(x, y, y') \right| \le M$

then the abovementioned BVP has a unique solution.

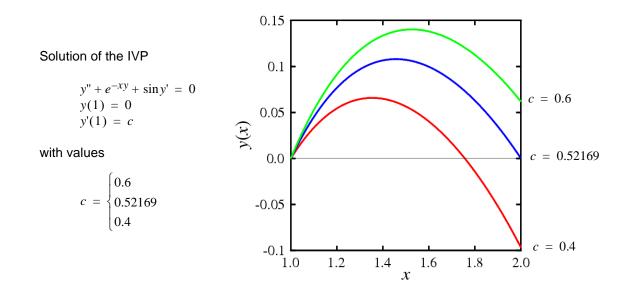
- An example:

$$y'' + e^{-xy} + \sin y' = 0, x \in [1, 2], y(1) = 0, y(2) = 0$$

or
$$y'' = f(x, y, y'), f(x, y, y') = -e^{-xy} - \sin y'.$$

Now
$$\frac{\partial}{\partial y} f(x, y, y') = xe^{-xy} > 0$$
, $\left| \frac{\partial}{\partial y'} f(x, y, y') \right| = \left| -\cos y' \right| \le 1$.

So, the equation has a unique solution.



Scientific computing III 2013: 13. Differential equations

Differential equations: boundary value problems

- The linear form of the 2nd order BVP can be written as

$$y'' = p(x)y' + q(x)y + r(x), a \le x \le b, y(a) = \alpha, y(b) = \beta.$$
 (1)

- For this form the existence of a unique solution requires that
 - a) p(x), q(x), r(x) are continuous in [a, b] and
 b) q(x) > 0 in [a, b].
- Consider two IVPs based on the BVP above:

 $\begin{aligned} y'' &= p(x)y' + q(x)y + r(x), \ x \in [a, b], \ y(a) &= \alpha, \ y'(a) = 0, \\ \text{solution } y_1(x) \\ y'' &= p(x)y' + q(x)y, \ x \in [a, b], \ y(a) = 0, \ y'(a) = 1, \end{aligned}$

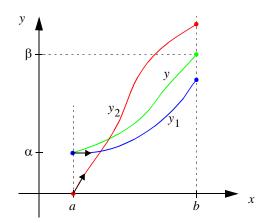
solution
$$y_2(x)$$

- It is easy to show that the unique solution to the original equation (1) can be expressed in terms of $y_1(x)$ and $y_2(x)$:

$$y(x) = y_1(x) + \frac{\beta - y_1(b)}{y_2(b)}y_2(x)$$

assuming $y_2(b) \neq 0^1$.

- A demo using GSL.



- In the **multiple shooting method** (**MSM**) the interval is divided into smaller parts and the SM is applied to every part so that the parts join to each other continuously.
 - Using the method you end up solving a (possibly nonlinear) group of equations.
 - Divide the interval [a, b] to parts

$$a = x_1 < x_2 < \dots < x_m = b$$
.

- Let $\mathbf{y}'(x;x_k,\mathbf{s}_k)$ be the solution to the IVP $\mathbf{y}' = \mathbf{f}(x,\mathbf{y})$, $\mathbf{y}(x_k) = \mathbf{s}_k$ on the interval $x \in [x_k, x_{k+1}]$.
- In the **MSM** method we pursue the vectors s_k , k = 1, 2, ..., m-1 that satisfy the continuity and boundary conditions

$$\mathbf{y}(x_{k+1};x_k,\mathbf{s}_k) = \mathbf{s}_{k+1}, \ k = 1, 2, ..., m-2$$

$$\mathbf{g}(\mathbf{s}_1, \mathbf{y}(b;x_{m-1},\mathbf{s}_{m-1})) = 0.$$

- The unknown initial values $\mathbf{s} = (\mathbf{s}_1 \mathbf{s}_2 \dots \mathbf{s}_{m-1})^T$ are obtained from the nonlinear system of equations

$$\mathbf{F}(\mathbf{s}) = \begin{bmatrix} \mathbf{y}(x_2; x_1, \mathbf{s}_1) - \mathbf{s}_2 \\ \mathbf{y}(x_3; x_2, \mathbf{s}_2) - \mathbf{s}_3 \\ \dots \\ \mathbf{g}(\mathbf{s}_1, \mathbf{y}(x_m; x_{m-1}, \mathbf{s}_{m-1})) \end{bmatrix} = 0$$

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Differential equations: boundary value problems

- The finite difference method is not based on solving the corresponding IVP but the BVP is discretized to a finite mesh and derivatives are approximated by finite differences.
 - An example shows how the method works¹
 - The BVP is

$$y'' = f(x, y, y')$$
$$y(a) = \alpha$$
$$y(b) = \beta.$$

- Divide interval [a, b] to m parts: $h = \frac{b-a}{m}$, $x_i = a + ih$, i = 0, ..., m
- Approximate the derivatives in the BVP by finite differences:

$$y'(x_n) \approx \frac{y_{n+1} - y_{n-1}}{2h}$$

 $y''(x_n) \approx \frac{y_{n+1} - 2y_n + y_{n-1}}{h^2}$

- Now we get the difference equation

$$y_{n+1} - 2y_n + y_{n-1} = h^2 f\left(x_n, y_n, \frac{y_{n+1} - y_{n-1}}{2h}\right), \quad n = 1, 2, ..., m-1$$

$$y_0 = \alpha, \ y_m = \beta.$$

^{1.} J.Haataja et al., Numeeriset menetelmät käytännössä, example 7.5.2

- In matrix form

$$\mathbf{A}\mathbf{y} = h^{2}\mathbf{f}(\mathbf{y})$$

$$\mathbf{A} = \begin{bmatrix} -2 & 1 & 0 & \dots & 0 & 0 \\ 1 & -2 & 1 & \dots & 0 & 0 \\ 0 & 1 & -2 & \dots & 0 & 0 \\ \dots & \dots & \dots & \dots & \dots \\ 0 & 0 & 0 & \dots & -2 & 1 \\ 0 & 0 & 0 & \dots & 1 & -2 \end{bmatrix}, \quad \mathbf{y} = \begin{bmatrix} y_{1} \\ y_{2} \\ y_{3} \\ \dots \\ y_{m-1} \end{bmatrix}, \quad \mathbf{f}(\mathbf{y}) = \begin{bmatrix} f\left(x_{1}, y_{1}, \frac{y_{2} - y_{0}}{2h}\right) - \alpha \\ f\left(x_{2}, y_{2}, \frac{y_{3} - y_{1}}{2h}\right) \\ f\left(x_{3}, y_{3}, \frac{y_{4} - y_{2}}{2h}\right) \\ \dots \\ f\left(x_{m-1}, y_{m-1}, \frac{y_{m} - y_{m-2}}{2h}\right) - \beta \end{bmatrix}$$

- Note that this is a nonlinear group of equations.

2

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ifferential equations: boundary value problems

- There are also methods based on expansion of the solution in terms of basis functions; e.g. collocation method (CM), Galerkin's method.
- Take for example the equation

$$y'' = f(x, y, y'), y(a) = \alpha, y(b) = \beta$$

- We expand the solution in terms of basis functions $\{\phi_1, \phi_2, ..., \phi_N\}$:

$$y(x) = \sum_{i=1}^{N} c_i \phi_i(x).$$

- The basis set can be chosen in an appropriate way considering the boundary conditions.
- In the **CM** method the above expansion is substituted in the equation and we require that it solves the equation in N points x_i within the interval [a, b]:

$$\sum_{i=1}^{N} c_i \phi_i''(x_j) = f\left(x_j, \sum_{i=1}^{N} c_i \phi_i(x_j), \sum_{i=1}^{N} c_i \phi_i'(x_j)\right), j = 1, 2, ..., N.$$

- From this equation we can determine the coefficients $\{c_1, c_2, ..., c_N\}$.

Practical tools:

Matlab: see help funfun for more information.

NAG: many routines for BVPs.

SLATEC: many routines for BVPs.

NETLIB: see the package ode (http://www.netlib.org/ode/)

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